Complex Analytic Methods

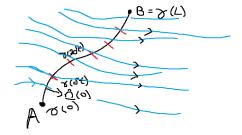
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Flux across a curve

Let $\gamma : [0, L) \to \mathbb{C}$ be a smooth regular curve. $|\gamma'(t)| = 1$ where $\gamma(t) = x(t) + iy(t)$



Let the velocity field be given by V(x, y) = u(x, y) + iv(x, y) and the unit normal vector $n(t) = \frac{\gamma'(t)}{i} = -i\gamma'(t)$ The flux across γ is approximated by $\frac{1}{N}\sum_{i=1}^{N} V(\gamma(\frac{L*j}{N})) \cdot n(\frac{L*j}{N})$

Flux across a curve

This tends to $\int_0^L V(\gamma(t)) \cdot n(t) dt$ as $N \to +\infty$ Easy to show that the flux equals $\int_0^L u(\gamma(t))y'(t) - v(\gamma(t))x'(t) dt$

Consider some differentiable $\psi : \mathbb{R}^2 \to \mathbb{R}$ $\frac{d}{dt}\psi(x(t), y(t)) = \frac{\partial\psi}{\partial x}x'(t) + \frac{\partial\psi}{\partial y}y'(t)$ Want this to equal -v(x(t), y(t))x'(t) + u(x(t), y(t))y'(t)

$$rac{\partial \psi}{\partial x} = -v(x,y)$$

 $rac{\partial \psi}{\partial y} = u(x,y)$

Want this so that flux across any curve joining any two points to be equal. This is required so that the solution makes sense physically (incompressibility)

When does such a ψ exist?

Choose some a in our domain.

 $\psi(x, y) = \int_{\Im(a)}^{y} u(\Re(a), \eta) \, d\eta - \int_{\Re(a)}^{x} v(\xi, y) \, d\xi$ Such a ψ satisfies the above equations. Remark: Note that for ψ as above to satisfy the system of equations, u and v need only be locally integrable.

We can start thinking about weak solutions to these equations...

Governing equation for ψ and vorticity

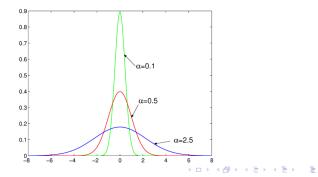
$$\begin{split} \frac{\partial^2 \psi}{\partial x^2} &+ \frac{\partial^2 \psi}{\partial y^2} = -\frac{\partial v}{\partial x} + \frac{\partial u}{\partial y} \\ \text{In fact, } \nabla \wedge V &= \left(\frac{\partial v}{\partial x} - \frac{\partial u}{\partial y}\right) \mathbf{k} \\ \omega &:= \left(\nabla \wedge V\right) \cdot \mathbf{k} \text{ (Vorticity)} \\ \Rightarrow \Delta \psi &= -\omega(x, y) \text{ (which is a Linear PDE)} \end{split}$$

Fundamental Solution

We can then ask what the Fundamental Solution is (use Theory of Distributions)

i.e. we seek
$$\psi \in D'(\mathbb{R}^2)$$
 s.t. $riangle \psi = \delta$
where $\delta \in D'(\mathbb{R}^2)$ is the Dirac Delta "function"

The Dirac Delta is not a function but you can think of it as a limit of smooth functions which go to $+\infty$ at 0 and tend to 0 everywhere else.



Fundamental Solution and Integral Representation of ψ

You can check that $riangle_{2\pi} \log |x| = \delta$ in \mathbb{R}^2 For relatively nice ω we

have that
$$\psi(x) = rac{1}{2\pi}\int_{\mathbb{R}^2}\omega(y)\log\lvert x-y
vert\,dy$$

Modelling Vortices

Using our previous interpretation of δ , we can try to model vortices as points of "infinite" singular vorticity. Suppose we have vortices at $a_1, ..., a_n$ of "strength" $c_1, ..., c_n$, then we model this as:

$$\omega = \sum_{k=1}^{n} c_k \delta(x - a_k)$$

where there is a proper way to make sense of $\delta(x - a_k)$

Then
$$\psi(x) = \sum_{k=1}^{n} c_k \log|x - a_k|$$

Notice that $\Delta \psi = 0$ except at $a_1, ..., a_n$

Laplace's equation and Boundary Conditions

So we focus on open sets where $\bigtriangleup\psi=0$ and direct our efforts towards solving this equation.

Some remarks on boundary conditions:

Through physical considerations one can derive the boundary condition $\psi={\rm constant}$ on a solid boundary.

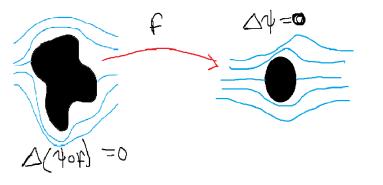
Curves on which $\psi={\rm constant}$ are called streamlines and are very important in visualizing flows.

When the flow is steady (doesn't change with time) streamlines are also particle paths.

Conformal Mapping Strategy

Strategy:

- Solve an equivalent flow outside the unit disk (Find ψ s.t. $\Delta \psi = 0$ outside unit disk)
- Find a conformal map (f'(z) ≠ 0) which maps desired region to the region outside the unit disk.
- Then $riangle \psi(f(z)) = 0$ in our desired region



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Conformal Mapping Strategy

Advantages:

- Gives exact solution
- Can be analyzed with more ease
- Can ensure the existence of f under certain conditions

Disadvantages:

 Practically impossible to efficiently use this method for complicated regions

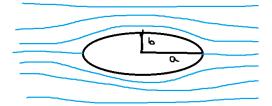
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Conformal Mapping Example

Want to find the flow past an ellipse, which is uniform (with speed U) away from ellipse

Solution for the unit disk is

$$\psi = U(y - \frac{y}{r^2}) = \Im(U(z + \frac{1}{z})) = \Im(w(z))$$



Finding f

We parametrise the boundary of the ellipse as

$$\gamma(t) = a\cos(t) + ib\sin(t) = \frac{a+b}{2}e^{it} + \frac{a-b}{2}e^{-it} (a > b)$$

Note that e^{it} parametrises the boundary of the unit disk. So the mapping $g(z) = \alpha_1 z + \frac{\alpha_2}{z}$ maps the boundary of the unit disk to the boundary of the ellipse in question. We want to invert this

$$g = \alpha_1 z + \frac{\alpha_2}{z}$$
 iff $\alpha_1 z^2 - gz + \alpha_2 = 0$ iff $z = \frac{g \pm \sqrt{g^2 - 4\alpha_1 \alpha_2}}{2\alpha_1}$

Turns out the correct sign to choose is +, i.e $f(z) = \frac{z + \sqrt{z^2 - 4\alpha_1 \alpha_2}}{2\alpha_1}$

Logarithmic Conjugation Theorem:

 Ω is a finitely connected region, with $K_1, ..., K_N$ the bounded components of its complement.

For each j, let a_j be a point in K_j .

If u is a real valued harmonic function on Ω then:

- there exist an analytic function f on Ω
- ▶ and there exist real numbers $c_1, ..., c_N$ such that

 $u(z) = \Re(f(z)) + c_1 \log|z - a_1| + ... + c_N \log|z - a_N|$

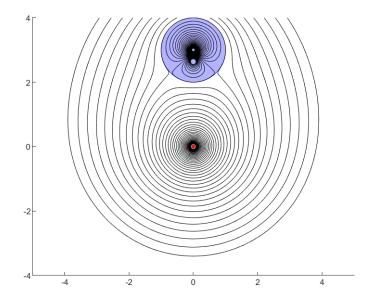
Runge Approximation Theorem: Suppose $K \subset \mathbb{C}$ is compact and f is analytic in KThen f can be uniformly approximated by rational functions with poles in $K^{\mathbb{C}}$ (rational functions approximate f in K)

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Combine these two and set up a linear system of equations to fit for the coefficients of the rational function using least squares

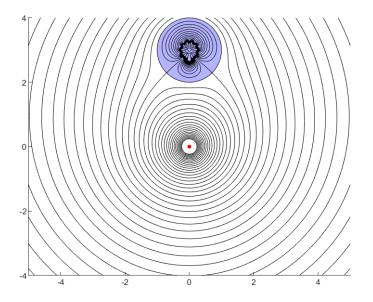
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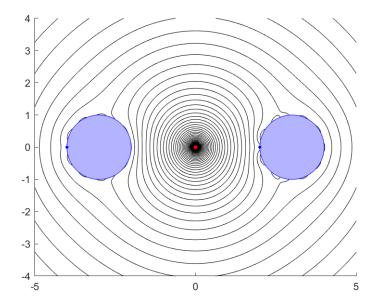
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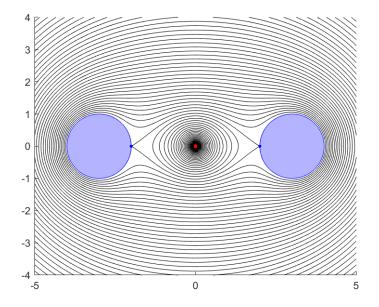
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Can indicate when boundary conditions are not achievable



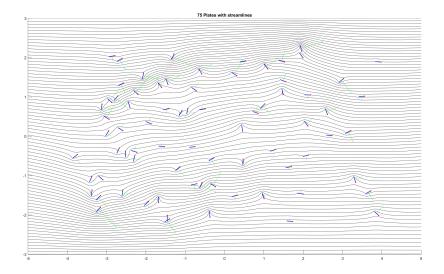
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Linear fit of boundary conditions



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75 plates in channel flow



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